

WIND INDUCED PRESSURES

Introduction

Both observation and experiment have shown that the atmospheric pressures experienced by surfaces exposed to the wind are different from the pressures on those same surfaces when the air is still but at the same density and temperature. Windward surfaces tend to have higher pressures, while leeward surfaces and surfaces parallel to the wind tend to have lower pressures.

This paper is an attempt to describe and explain wind-induced pressure changes using statistical mechanics and kinetic gas theory.

Pressures in the Free Atmosphere

In this paper (as in most of the papers in this collection), we do not treat air pressure as an intrinsic property of the mixture of gases that make up our atmosphere. Instead, we consider it to be generally a function of both scale of place and scale of time, and specifically a function of the relationship between the air molecules and the sensor.

In dealing with parcels of air, atmospheric scientists are accustomed to the fact that the pressures sensed in one part of the parcel will be different from those sensed in a different location. Moreover, we normally expect that these same pressures will change significantly with time as the parcel moves across the landscape. That is why our symbol for pressure (\bar{p}_i) has a bar over it. This bar denotes that the term is the mean of several values, and not a single individual value. The i subscript denotes that pressure is a parameter of the interactive sub-population of molecules. The parameter only exists when the molecules of the atmosphere interact with an object of interest.

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We also recognize that the pressure sensed by a sensor will have one value when both the air and the sensor are still, other values (depending upon wind direction and velocity relative to the sensor) when the wind is blowing, a third set of pressure values when the sensor is moving in still air, and a fourth set of values when both the sensor is moving and the wind is blowing.

No. Atmospheric pressures are not simple things.

Consequently, the reader is strongly urged to be familiar with the concepts and equations in both [The Nature of Atmospheric Pressures](#) and in [The Nature of Wind](#). Without that familiarity, most of what follows will make little sense.

With that caution in mind, let us proceed to recapitulate the contents of those two papers.

The Nature of Atmospheric Pressures

As we saw in [The Nature of Atmospheric Pressures](#), the atmospheric pressure on any surface is the product of the mean number of molecules striking that surface in unit time (the mean molecular flux) and the mean impulse that these impacts transfer to that surface. This is illustrated by the *Universal Pressure Equation*:

$$\bar{p}_i = \bar{f}_i \bar{\Phi}_i \quad \text{WIP01}$$

Here, \bar{p}_i is the mean pressure per square meter in Pascals, \bar{f}_i is the mean molecular flux in number of molecular impacts per square meter per second, and $\bar{\Phi}_i$ is the mean molecular impulse per impact in newtons.

Being essentially a definition, this universal pressure equation is valid for ideal gases and for real gases, under conditions of equilibrium and under conditions of non-equilibrium, in the laboratory and in the free atmosphere, for still air and for moving air. It is even valid for moving objects in moving air.

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Limitations of the One-Dimensional Notation System – Our current notational system is a [*one-dimensional notation system*](#) with p and i subscripts to denote **molecular** directions of movement and speeds relative to the object of interest. We need a second notation system to denote **wind** direction and speed relative to our object of interest, which in this paper will be the sensing surface of an imaginary sensor.

To do this, we will postulate a standard tri-axial orthogonal reference system; that is, an x -axis, a y -axis and a z -axis—with all three axes at right angles to one another. This system does not replace our p -axis, it supplements it. Winds are three-dimensional phenomena, therefore we need a three-dimensional notation system to describe them. This same notational system is used in other papers in this collection that also refer to winds.

Wind System Parameters – In our discussion of winds, we will make use of the following parameters. Unless these terms bear a directional superscript or subscript as defined below, they are solely scalar terms. When they do bear such a superscript or subscript, the parent terms are the scalar component of a vector expression.

$\bar{\omega}$ Mean wind speed along the wind axis in meters per second.

a Error function of the mean wind speed expressed in standard deviations: $a = \text{erf}(\bar{\omega} / \sigma)$. Note that $a + b = 0.5$.

b Complimentary error function of the mean wind speed expressed in standard deviations: $b = \text{erfc}(\bar{\omega} / \sigma)$. Note that $a + b = 0.5$.

α Angle of incidence of the wind vector to the sensing surface in degrees. When the wind is normal to and toward that surface, α is $+90^\circ$. When the wind is parallel to that surface, α is 0° . When the wind is normal to and away from that surface, α is -90° . Intermediate angles of incidence will have intermediate values of α . These values will range from $+90^\circ$ through 0° to -90° .

\bar{T}_i Mean temperature sensed by the sensing surface, in Kelvins.

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- \bar{p}_i Mean pressure in Pascals, measured normal to the sensor.
- \bar{n} Mean molecular number density in number of molecules per cubic meter.
- \bar{m}_i Mean molecular impulse mass in kilograms. This is not the same as the mean molecular mass of the population of molecules. See [Molecular Masses](#) for the difference.
- \bar{u}_p Mean component molecular kinetic energy of translation in joules, measured normal to and toward the sensor along a single axial arm.
- Σu_p The sum of those molecular kinetic energies of translation in joules for a particular axial arm—in this case, the *p* arm.
- \bar{v} Mean molecular speed in meters per second. Subscripts will indicate whether the term refers to the general population (*p*) or the interactive sub-population (*i*).
- k_B Boltzmann's Constant in joules per molecule per Kelvin.
- \bar{f}_i Mean molecular flux in number of molecules per square meter per second, measured normal to and toward the sensor.
- $\bar{\phi}_i$ Mean molecular impulse per impact in joules, measured normal to and toward the sensor.
- σ *Sigma*. This is the axial root-mean-square molecular speed in meters per second in still air at temperature \bar{T} and mean molecular impulse mass \bar{m}_i . As such, it is also the standard deviation (σ) of the axial molecular speed distribution in still air. As a speed and not a velocity, it is a scalar term, possessing only magnitude.

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Superscripts and subscripts: The following terms may be used as either subscripts or as superscripts, depending upon which position is available.

p Denotes membership in that portion of the general population of molecules that has a component of movement normal to and toward the object of interest.

$-p$ Denotes membership in that portion of the general population of molecules that has a component of movement normal to and away from the object of interest.

i Denotes membership in that portion of the interactive sub-population of molecules whose next or current interaction is with the object of interest.

\rightarrow Denotes that the molecules designated by the parent expression have a component of motion in the windward direction.

\leftarrow Denotes that the molecules designated by the parent expression have a component of motion in the leeward direction.

\downarrow Denotes that the molecules designated by the parent expression have a component of motion normal to and toward the wind axis.

The Nature of Wind

We were able to capture pressure with a single equation, WIP01. Wind is a far more elusive quarry. From the standpoint of kinetic gas theory and statistical mechanics, we do not, as yet, have a single equation that captures both the wind direction and its force or magnitude. [The Nature of Wind](#) gives us several such equations, depending upon the angle of incidence.

Nevertheless, let us continue our discussion of wind induced pressures by defining once more just exactly what we mean by wind:

When a wind is blowing, a stationary reference system will measure more wind molecules with a component of movement in the windward

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direction than molecules having a component of movement in any other direction. In addition, this system will measure the windward molecules as moving with a higher mean speed than the mean molecular component speed in any other direction.

A reference system moving with the wind will measure the same molecular parameters as it would in still air at the same temperature, pressure, and molecular number density. The terms and expressions in the equation of state:

$$\bar{p}_i = \bar{n}k_B\bar{T} = \bar{n}\bar{m}_i\bar{\sigma}^2 \quad \text{WIP02}$$

have exactly the same numerical values in both still air and when the wind is blowing—as long as the sensors are moving with the wind.

As we will see below, the numerical values of these parameters (except for k_B and \bar{m}_i) will change significantly when measured relative to a stationary reference system when the wind blows.

Variability of Pressures when the Wind Blows

Both observation and experiment have shown that pressure sensors in moving air change their readings in response to changes in both wind velocity and in the angle of incidence that the wind makes with the sensing surface.

Consequently, we will examine this variability on a windward sensing surface, a leeward sensing surface, and a sensing surface that is parallel to the wind direction. From these results, we may calculate the expected values on sensing surfaces at intermediate angles of incidence.

Pressure Variability on a Windward Surface

Let us postulate a wind blowing through an imaginary frame of reference. Table 4 in [The Nature of Winds](#) gives us the following increment to the static pressure upon a windward sensor:

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$$\Delta\bar{p}_i = \bar{n}\bar{m}_i\bar{\omega}^2 + a\bar{n}\bar{m}_i\bar{\omega}^2 \quad \text{WIP03}$$

This results in a total dynamic pressure on the windward surface of:

$$\bar{p}_i^{\rightarrow} = \bar{p}_i + \left(\bar{n}\bar{m}_i\bar{\omega}^2 + a\bar{n}\bar{m}_i\bar{\omega}^2 \right) \quad \text{WIP04}$$

Pressure Variability on a Leeward Surface

Table 5 in [The Nature of Winds](#) gives us the following decrement to the static pressure upon a leeward sensor:

$$\Delta\bar{p}_i = -2b\bar{n}\bar{m}_i\bar{\omega}^2 \quad \text{WIP05}$$

This results in a total dynamic pressure on the leeward surface of:

$$\bar{p}_i^{\leftarrow} = \bar{p}_i - 2b\bar{n}\bar{m}_i\bar{\omega}^2 \quad \text{WIP06}$$

WARNING: These leeward pressures should be viewed strictly as limiting cases. In the free atmosphere, surfaces to the leeward of obstructions to the wind will virtually always be subject to turbulence of one sort or another. Consequently, both spatial and temporal variations in leeward pressures having significant magnitude are the rule.

Pressure Variability on a Parallel Surface

Table 6 in [The Nature of Winds](#) gives us the following decrement to the static pressure upon a sensor that is parallel to the wind axis:

$$\Delta\bar{p}_i = -a\bar{n}\bar{m}_i\bar{\omega}^2 \quad \text{WIP07}$$

This results in a total dynamic pressure on the parallel surface of:

$$\bar{p}_i^{\downarrow} = \bar{p}_i - a\bar{n}\bar{m}_i\bar{\omega}^2 \quad \text{WIP08}$$

Pressure Variability on Intermediate Surfaces

The Nature of Winds gives us the procedure for calculating the resulting incident pressures when the angle of incidence of the wind lies somewhere between normal to and toward the sensor and parallel to the sensor. That procedure is:

$$\bar{p}_i^\alpha = (\sin \alpha) \bar{p}_i^\rightarrow + (1 - \sin \alpha) \bar{p}_i^\downarrow \quad \text{WIP09}$$

The Nature of Winds also gives us the procedure for calculating the resulting incident pressures when the angle of incidence of the wind lies somewhere between normal to and away from the sensor and parallel to the sensor. That procedure is:

$$\bar{p}_i^\alpha = (\sin \alpha) \bar{p}_i^\leftarrow + (1 - \sin \alpha) \bar{p}_i^\downarrow \quad \text{WIP10}$$

TABLES

Introduction: The tables in this section give numerical values for selected atmospheric parameters associated with wind-induced pressures at selected wind speeds and selected angles of incidence.

In order to do this, we have to start with a standard set of still air parameters and conditions. These initial conditions include both an absence of wind and an absence of any changes of state in the water vapor content of the atmosphere. Moreover, when the wind blows, all tables assume a completely laminar non-viscous flow, with no turbulence.

Still Air Parameters:

\bar{p}_i 1,000 hectopascals or 10^5 Pascals. This is the mean ambient atmosphere pressure in still air under conditions approaching equilibrium.

\bar{n} 2.4293×10^{25} molecules per cubic meter. This is the molecular number density that produces an ambient pressure of 1000 hectopascals in still air at a temperature of 25°C .

k_B $1.3806488 \times 10^{-23}$ joules per molecule per Kelvin. This is Boltzmann's Constant. It is the constant of proportionality between molecular kinetic energy of translation and temperature.

\bar{T}_i 25°C or 298.15°K . This is the ambient mean atmospheric temperature.

\bar{m}_i 4.777108×10^{-26} kilograms. This is the mean molecular impulse mass at a vapor pressure of 10^3 Pascals. At the specified temperature of 25°C , this corresponds to a relative humidity of about 32%.

σ 293.55 meters per second. This is the square-root of the mean of the squares of the deviations of the individual molecule speeds from the mean speed (0 meters per second) along any single axis of movement in still air. In kinetic gas

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theory, this speed is known as the root-mean-square speed (v^{rms}). In statistical mechanics, this speed is known as the standard deviation of the molecular axial speed distribution.

Governing Equation: In still air, all of the above parameters are related to one another by the governing equation:

$$\bar{p}_i = \bar{n}k_B\bar{T}_i = \bar{n}\bar{m}_i\sigma^2 \qquad \text{WIP11}$$

Properly speaking, this equation applies only to an ideal gas under conditions of equilibrium. However, both observation and experiment have shown that it can be applied to the mixture of real gases that make the free atmosphere under conditions approaching equilibrium. That is to say, that the real atmospheric values equal the theoretical values to within the normal limits of scientific precision.

In keeping with this principle, the reader should treat the very precise numbers in the following tables with a grain of salt. They simply represent the computer doing what it does, regardless of the scientific merits of the process. You can probably accept the accuracy of three significant figures in most cases or four—if conditions are really controlled. Otherwise, use as many significant figures as you feel comfortable with in estimating the accuracy of the wind speed.

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TABLE WIP01: PRESSURE CHANGES AT SELECTED WIND SPEEDS

Wind Speed	Mean Windward Pressure Increment	Mean Leeward Pressure Decrement	Mean Parallel Pressure Decrement
$\bar{\omega}$	$\Delta\bar{p}_i^{\rightarrow}$	$\Delta\bar{p}_i^{\leftarrow}$	$\Delta\bar{p}_i^{\downarrow}$
m sec ⁻¹	hPa	hPa	hPa
0	0.000	0.000	0.000
1	0.012	-0.012	-0.000
2	0.047	-0.046	-0.000
3	0.105	-0.104	-0.000
4	0.188	-0.184	-0.001
5	0.294	-0.286	-0.002
6	0.425	-0.411	-0.003
7	0.579	-0.558	-0.005
8	0.759	-0.727	-0.008
9	0.963	-0.917	-0.011
10	1.192	-1.129	-0.016
11	1.446	-1.362	-0.021
12	1.726	-1.617	-0.027
13	2.031	-1.892	-0.035
14	2.361	-2.188	-0.043
15	2.718	-2.505	-0.053
16	3.100	-2.842	-0.065
17	3.509	-3.199	-0.077
18	3.944	-3.576	-0.092
19	4.406	-3.973	-0.108
20	4.894	-4.390	-0.126

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TABLE WIP02: PRESSURE CHANGES AT SELECTED WIND SPEEDS

Wind Speed	Mean Windward Pressure Increment	Mean Leeward Pressure Decrement	Mean Parallel Pressure Decrement
$\bar{\omega}$	$\Delta\bar{p}_i^{\rightarrow}$	$\Delta\bar{p}_i^{\leftarrow}$	$\Delta\bar{p}_i^{\downarrow}$
m sec ⁻¹	hPa	hPa	hPa
0	0.000	0.000	0.000
5	0.294	-0.286	-0.002
10	1.192	-1.129	-0.016
15	2.718	-2.505	-0.053
20	4.894	-4.390	-0.126
25	7.745	-6.761	-0.246
30	11.295	-9.594	-0.425
35	15.565	-12.867	-0.675
40	20.581	-16.556	-1.006
45	26.363	-20.637	-1.432
50	32.937	-25.089	-1.962
55	40.323	-29.888	-2.609
60	48.544	-35.012	-3.383
65	57.624	-40.439	-4.296
70	67.583	-46.147	-5.359
75	78.442	-52.114	-6.582
80	90.225	-58.320	-7.976
85	102.951	-64.742	-9.552
90	116.641	-71.361	-11.320
95	131.315	-78.156	-13.290
100	146.994	-85.107	-15.472

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TABLE WIP03: PRESSURE CHANGES AT SELECTED ANGLES OF INCIDENCE AND SELECTED WIND SPEEDS

Angle of Incidence	Change at Wind Speed of	Change at Wind Speed of	Change at Wind Speed of	Change at Wind Speed of
α	5	10	15	20
$^{\circ}$	ms^{-1}	ms^{-1}	ms^{-1}	ms^{-1}
90	0.294	0.294	0.294	0.294
85	0.293	0.293	0.293	0.293
80	0.289	0.289	0.289	0.289
75	0.284	0.284	0.284	0.284
70	0.276	0.276	0.276	0.276
65	0.266	0.266	0.266	0.266
60	0.254	0.254	0.254	0.254
55	0.240	0.240	0.240	0.240
50	0.225	0.225	0.225	0.225
45	0.207	0.207	0.207	0.207
40	0.188	0.188	0.188	0.188
35	0.168	0.168	0.168	0.168
30	0.146	0.146	0.146	0.146
25	0.123	0.123	0.123	0.123
20	0.099	0.099	0.099	0.099
15	0.075	0.075	0.075	0.075
10	0.049	0.049	0.049	0.049
5	0.024	0.024	0.024	0.024
0	-0.002	-0.002	-0.002	-0.002

All changes are in hectopascals.

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TABLE WIP04: PRESSURE CHANGES AT SELECTED ANGLES OF INCIDENCE AND SELECTED WIND SPEEDS

Angle of Incidence	Change at Wind Speed of	Change at Wind Speed of	Change at Wind Speed of	Change at Wind Speed of
α	25	30	35	40
$^{\circ}$	ms ⁻¹	ms ⁻¹	ms ⁻¹	ms ⁻¹
90	7.735	11.280	15.545	20.553
85	7.705	11.235	15.483	20.471
80	7.614	11.102	15.299	20.226
75	7.463	10.881	14.992	19.819
70	7.254	10.574	14.567	19.253
65	6.988	10.183	14.025	18.534
60	6.666	9.712	13.372	17.665
55	6.292	9.163	12.612	16.655
50	5.868	8.542	11.751	15.510
45	5.398	7.852	10.795	14.239
40	4.884	7.099	9.752	12.853
35	4.332	6.289	8.629	11.361
30	3.745	5.428	7.436	9.775
25	3.127	4.522	6.181	8.106
20	2.484	3.579	4.874	6.369
15	1.820	2.605	3.524	4.575
10	1.140	1.608	2.143	2.739
5	0.450	0.596	0.740	0.875
0	-0.246	-0.424	-0.673	-1.004

All changes are in hectopascals.

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TABLE WIP05: PRESSURE CHANGES AT SELECTED ANGLES OF INCIDENCE AND SELECTED WIND SPEEDS

Angle of Incidence	Change at Wind Speed of	Change at Wind Speed of	Change at Wind Speed of	Change at Wind Speed of
α	45	50	55	60
$^{\circ}$	ms ⁻¹	ms ⁻¹	ms ⁻¹	ms ⁻¹
90	26.328	32.893	40.269	48.479
85	26.223	32.760	40.106	48.282
80	25.907	32.363	39.617	47.691
75	25.383	31.705	38.808	46.712
70	24.654	30.791	37.683	45.352
65	23.728	29.627	36.252	43.621
60	22.610	28.223	34.525	41.532
55	21.309	26.590	32.515	39.101
50	19.834	24.739	30.238	36.347
45	18.198	22.685	27.712	33.291
40	16.413	20.443	24.954	29.955
35	14.492	18.031	21.987	26.366
30	12.450	15.467	18.832	22.551
25	10.302	12.770	15.515	18.538
20	8.065	9.961	12.059	14.359
15	5.755	7.062	8.492	10.044
10	3.391	4.094	4.841	5.628
5	0.990	1.079	1.133	1.143
0	-1.429	-1.958	-2.604	-3.377

All changes are in hectopascals.

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TABLE WIP06: PRESSURE CHANGES AT SELECTED ANGLES OF INCIDENCE AND SELECTED WIND SPEEDS

Angle of Incidence	Change at Wind Speed of	Change at Wind Speed of	Change at Wind Speed of	Change at Wind Speed of
α	65	70	75	80
$^{\circ}$	ms ⁻¹	ms ⁻¹	ms ⁻¹	ms ⁻¹
90	57.546	67.491	78.335	90.101
85	57.310	67.213	78.012	89.728
80	56.606	66.384	77.045	88.611
75	55.439	65.009	75.442	86.760
70	53.817	63.098	73.215	84.187
65	51.752	60.666	70.380	80.914
60	49.262	57.732	66.960	76.963
55	46.363	54.318	62.980	72.367
50	43.079	50.449	58.471	67.159
45	39.435	46.156	53.467	61.379
40	35.458	41.471	48.006	55.072
35	31.178	36.430	42.130	48.285
30	26.629	31.071	35.883	41.070
25	21.844	25.434	29.313	33.482
20	16.860	19.564	22.470	25.578
15	11.716	13.503	15.405	17.419
10	6.449	7.300	8.174	9.067
5	1.101	1.000	0.830	0.585
0	-4.288	-5.349	-6.570	-7.961

All changes are in hectopascals.

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The Luffing Angle: The angle at which the windward force becomes a leeward force is known as the luffing angle. In sailing, this luffing will always occur within five degrees of the wind. This is shown in the tables above. Note that the higher the wind velocity, the closer to five degrees that the luffing angle becomes. Given the extreme variability of both the wind direction and the wind speed under natural conditions, a more precise value than this is ill-advised.

REFERENCES

INTERNAL REFERENCES: These are other papers in this collection that are cited or linked during the course of the discussion.

[The Nature of Atmospheric Pressures](#) – This paper defines gas pressures in terms of kinetic gas theory and statistical mechanics.

[The Nature of Winds](#) – This paper defines winds in terms of kinetic gas theory and statistical mechanics.

[Molecular Flows](#) – This paper gives the derivation of the Universal Molecular Flux equation.

[Molecular Impulses](#) – This paper gives the derivation of the Universal Molecular Impulse equation.

[The Equipartition Conundrum](#) – This paper shows how and why the intermolecular transfer of momentum is virtually solely translational.

[Molecular Speeds and Velocities](#) – This paper shows how the various measures of molecular velocities are derived and their mathematical relationships.

[Kinetic Energies of Translation](#) – This paper shows how the various measures of molecular kinetic energies of translation are derived and their mathematical relationships to one another and to the thermal term $k_B \bar{T}$.

[Molecular Masses](#) – This paper shows how the values for the various atmospheric molecular masses are calculated.

[The Probability Density Curve](#) – This paper gives a brief treatment of the history and utility of the probability density curve.

[The Balloon and the Tower](#) – This is a thought experiment involving a drifting balloon and an airfoil-shaped tower. It shows how orientation to the wind affects both temperatures and pressures.

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EXTERNAL REFERENCES: These are papers by other authors that contain specific statements or bits of data that are specifically incorporated in the above discussion. Specifically, these papers treat gaseous flows from the standpoints of kinetic gas theory and statistical mechanical theory.

D. Tabor; **Gases, Liquids, and Solids**; Third Edition; Cambridge University Press, 1991.

Arthur Brown; **Statistical Physics**; Elsevier, New York, 1970.

James Jeans; **An Introduction to the Kinetic Theory of Gases**; Cambridge Library Collection, 1940.

GENERAL REFERENCES: These are papers by other authors that contain general treatments of kinetic gas theory, statistical mechanics and thermodynamics, atmospheric physics, and other scientific fields that are used in the above discussion.

Charles Kittel; **Thermal Physics**; John Wiley & Sons, New York, 1969.

Wolfgang Pauli; **Statistical Mechanics**; Dover Press, Mineola, 1973.

William D. Sellers, **Physical Climatology**; University of Chicago Press, Chicago, 1965.

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Richmond W. Longley; **Elements of Meteorology**, John Wiley & Sons, Inc.; New York, 1970.

Herber Riehl; **Introduction to the Atmosphere**; McGraw-Hill Book Company, New York, 1972.